

Asymptotic spectral analysis of Toeplitz operators on symplectic manifolds

Yuri A. Kordyukov

Ufa Federal Research Centre RAS and Novosibirsk State University

Topological structures in mathematics, physics and biology,
Novosibirsk, September, 2018

The setting

X a (compact) smooth manifold;

(L, h^L, ∇^L) a Hermitian line bundle on X :

- $L \rightarrow X$ a complex line bundle on X :
locally, over some open $\Omega \subset X$,
 $L|_{\Omega} \cong \Omega \times \mathbb{C}$; $C^\infty(\Omega, L|_{\Omega}) \cong C^\infty(\Omega)$;
- h^L a Hermitian structure in the fibers of L :

$$s, s' \in L \rightarrow (s, s')_{h^L} \in \mathbb{C},$$

- ∇^L a connection (covariant derivative): for $U \in C^\infty(X, TX)$

$$\nabla_U^L : C^\infty(X, L) \rightarrow C^\infty(X, L),$$

which is Hermitian:

$$\nabla_U^L (s, s')_{h^L} = (\nabla_U^L s, s')_{h^L} + (s, \nabla_U^L s')_{h^L}, \quad s, s' \in C^\infty(X, L).$$

Example

- $X = \mathbb{R}^d$.
- $L = X \times \mathbb{C} \rightarrow X$ the trivial line bundle, $C^\infty(X, L) \cong C^\infty(X)$.
- The Hermitian structure is given by $h \in C^\infty(X)$: for $Z \in \mathbb{R}^d$

$$|s|_h^2 = h(Z)|s|^2, \quad s \in L_Z = \{Z\} \times \mathbb{C};$$

- The connection: for any $U \in C^\infty(X, TX)$, $\nabla_U^L : C^\infty(X, L) \rightarrow C^\infty(X, L)$ is the first order differential operator:

$$\nabla_U^L = \frac{\partial}{\partial U} + \Gamma(U),$$

$\Gamma = \sum_{j=1}^d \Gamma_j(Z) dZ^j \in \Omega^1(X)$ is the connection one-form;

- ∇^L is Hermitian $\Leftrightarrow \Gamma + \bar{\Gamma} = -h^{-1} dh$.

The Bochner-Laplacian

Let (L, h^L, ∇^L) be a Hermitian line bundle on X .

The connection can be considered as an operator

$$\nabla^L = d + \Gamma : C^\infty(X, L) \rightarrow C^\infty(X, T^*X \otimes L)$$

Fix a Riemannian metric g on X .

We have L^2 -inner products on $C^\infty(X, L)$ and $C^\infty(X, T^*X \otimes L)$:

$$(s, s')_{L^2(X, L)} = \int_X (s(x), s'(x))_{h^L} dv_g(x), \quad s, s' \in C^\infty(X, L).$$

The formally adjoint operator

$$(\nabla^L)^* : C^\infty(X, T^*X \otimes L) \rightarrow C^\infty(X, L).$$

For $s \in C^\infty(X, L)$, $s' \in C^\infty(X, T^*X \otimes L)$:

$$(\nabla^L s, s')_{L^2(X, T^*X \otimes L)} = (s, (\nabla^L)^* s')_{L^2(X, L)}.$$

The Bochner-Laplacian

Definition

The **Bochner-Laplacian** Δ^L associated with a Hermitian line bundle (L, h^L, ∇^L) :

$$\Delta^L = (\nabla^L)^* \nabla^L : C^\infty(X, L) \rightarrow C^\infty(X, L).$$

If $\{e_j\}_{j=1, \dots, d}$ is a local orthonormal frame of TX , then

$$\Delta^L = - \sum_{j=1}^d \left[(\nabla_{e_j}^L)^2 - \nabla_{\nabla_{e_j}^{TX} e_j}^L \right],$$

where ∇^{TX} the Levi-Civita connection of g .

Example: magnetic Laplacian

- $X = \mathbb{R}^d$.
- $L = X \times \mathbb{C}$ the trivial line bundle.
- The Hermitian structure ($h(Z) = 1$):

$$|s(Z)|_h^2 = |s(Z)|^2.$$

- The connection form

$$\Gamma = -i\mathbf{A}, \quad \mathbf{A} = \sum_{j=1}^d A_j(Z) dZ_j$$

is a real-valued one form (a magnetic potential).

- The Riemannian metric is the standard metric $g = \sum_{j=1}^d dZ_j^2$.
- The Bochner-Laplacian is the magnetic Schrödinger operator:

$$\Delta^L = - \sum_{j=1}^d \left(\frac{\partial}{\partial Z_j} - iA_j(Z) \right)^2.$$

The curvature

Let (L, h^L, ∇^L) be a Hermitian line bundle on X .

The **curvature** of ∇^L is the differential two-form R^L on X :

$$R^L(U, V) = \nabla_U^L \nabla_V^L - \nabla_V^L \nabla_U^L - \nabla_{[U, V]}^L, \quad U, V \in TX.$$

For the connection $\nabla_U^L = \frac{\partial}{\partial U} + \Gamma(U)$, its curvature is given by

$$R^L = d\Gamma.$$

For the magnetic Laplacian $\Delta^L = -\sum_{j=1}^d \left(\frac{\partial}{\partial Z_j} - iA_j(Z) \right)^2$,

$$R^L = -i\mathbf{B}$$

where $\mathbf{B} = d\mathbf{A}$ is a real-valued two form (the magnetic field):

$$\mathbf{B} = \sum_{j, k=1}^d B_{jk}(Z) dZ_j \wedge dZ_k, \quad B_{jk} = \frac{\partial A_k}{\partial Z_j} - \frac{\partial A_j}{\partial Z_k}.$$

Non-degeneracy assumption

We will assume that the differential two form

$$\omega = \frac{i}{2\pi} R^L$$

is nondegenerate (of full rank):

$$(\omega(U, V) = 0 \text{ for any } V \in TX) \Rightarrow U = 0.$$

For the magnetic Laplacian $\Delta^L = -\sum_{j=1}^d \left(\frac{\partial}{\partial Z_j} - iA_j(Z) \right)^2$,

$$\omega = \frac{1}{2\pi} \mathbf{B}$$

is non-degenerate.

In particular, the dimension d is even: $d = 2n$.

Relation with geometric quantization

- (X, ω) a symplectic manifold, $\dim X = 2n$, so it is a classical phase space.
- A Hermitian line bundle (L, h^L, ∇^L) on X , satisfying:

$$\frac{i}{2\pi} R^L = \omega$$

is called a **prequantum line bundle**.

- (X, ω) is called **quantizable** \Leftrightarrow there exists a prequantum bundle ($\Leftrightarrow [\omega] \in H^2(X, \mathbb{Z})$).
- In geometric quantization scheme (Kostant-Souriau), the operators act on sections of L (prequantization).

Example: the 2-sphere

- X the two-dimensional sphere

$$S^2 = \{(x, y, z) \in \mathbb{R}^3 : x^2 + y^2 + z^2 = 1\},$$

equipped with the Riemannian metric induced by the standard Euclidean metric in \mathbb{R}^3 .

- ω is a scalar multiple of the volume 2-form dx_g :

$$\omega = s dx_g, \quad s \in \mathbb{R}.$$

- (X, ω) is quantizable \Leftrightarrow the area $4\pi s = n \in \mathbb{Z}$.
- The corresponding prequantum line bundle (L_n, ∇_n) is a well-known Wu-Yang magnetic monopole, which provides a natural topological interpretation of Dirac's monopole of magnetic charge $g = nh/2e$.

The renormalized Bochner-Laplacian

- $J_0 : TX \rightarrow TX$ a skew-adjoint linear endomorphism:

$$\omega(u, v) = g(J_0 u, v), \quad u, v \in TX;$$

- τ is a smooth function on X given by

$$\tau(x) = \pi \operatorname{Tr}[(-J_0^2(x))^{1/2}], \quad x \in X.$$

- L^p the p -th tensor power of L , $p \in \mathbb{N}$;
- $\nabla_U^{L^p} : C^\infty(X, L^p) \rightarrow C^\infty(X, L^p)$ the induced connection on L^p :

$$\nabla_U^{L^p} = \frac{\partial}{\partial U} + p\Gamma^L(U), \quad U \in TX.$$

Definition (V. Guillemin - A. Uribe, 1988)

The renormalized Bochner-Laplacian Δ_p acts on $C^\infty(X, L^p)$:

$$\Delta_p = \Delta^{L^p} - p\tau.$$

Magnetic Laplacian

- $X = \mathbb{R}^{2n}$, $L = X \times \mathbb{C}$ the trivial line bundle.
- The Hermitian structure $|s(Z)|_h^2 = |s(Z)|^2$.
- The connection form $\Gamma = -i\mathbf{A}$, where $\mathbf{A} = \sum_{j=1}^{2n} A_j(Z) dZ_j$ is a real-valued one form.
- The Bochner-Laplacian

$$\Delta^{L^p} = - \sum_{j=1}^{2n} \left(\frac{\partial}{\partial Z_j} - ipA_j(Z) \right)^2, \quad p = \frac{1}{\hbar}.$$

- $J_0 = \frac{1}{2\pi} B$, where $B : TX \rightarrow TX$ be a skew-adjoint operator

$$\mathbf{B}(u, v) = g(Bu, v), \quad u, v \in TX.$$

- $\tau(Z) = \frac{1}{2} \text{Tr}(B^* B)^{1/2} = \text{Tr}^+(B)$.

Complex manifolds

- $X = \mathbb{C}^n$, $L = X \times \mathbb{C}$ the trivial line bundle.
- The Hermitian structure is given by $h \in C^\infty(X)$: for $z = x + iy \in \mathbb{C}^n$

$$|s|_h^2 = h(z)|s|^2, \quad s \in L_z;$$

- The Hermitian connection

$$\nabla^L = d + \Gamma, \quad \Gamma + \bar{\Gamma} = -h^{-1}dh;$$

Assume that Γ is compatible with the complex structure of \mathbb{C}^n (a holomorphic Hermitian connection — the Chern connection), then, Γ is a $(1, 0)$ -form:

$$\Gamma = \partial \log h = \sum_{j=1}^n h^{-1} \frac{\partial h}{\partial z_j} dz_j;$$

Complex manifolds

- The curvature $R = d\Gamma$ is a purely imaginary 2-form: $(1, 1)$ -form

$$R = \bar{\partial}\partial \log h.$$

- For the symplectic form ω , we have

$$\omega = \frac{i}{2\pi} \bar{\partial}\partial \log h.$$

- ω is positive if $h = e^{-\varphi}$, $\varphi : X \rightarrow \mathbb{C}$ a smooth strictly plurisubharmonic function:

$$\omega = \frac{i}{2\pi} \sum_{j,k=1}^n \frac{\partial^2 \varphi}{\partial z_j \partial \bar{z}_k} dz_j \wedge d\bar{z}_k, \quad \left(\frac{\partial^2 \varphi}{\partial z_j \partial \bar{z}_k} \right)_{j,k=1,\dots,n} > 0.$$

Kähler manifolds

- A particular case: the Hermitian structure is given by

$$|s|_h^2 = h(z)|s|^2, \quad h(z) = e^{-\frac{\pi}{2}|z|^2};$$

- The connection form

$$\Gamma = \partial \log h = -\pi \sum_{j=1}^n \bar{z}_j dz_j;$$

- The symplectic form ω is the canonical symplectic form:

$$\omega = \frac{i}{2\pi} \bar{\partial} \partial \log h = \frac{i}{2} \sum_{j=1}^n dz_j \wedge d\bar{z}_j = \sum_{j=1}^n dx_j \wedge dy_j.$$

- J_0 is a complex structure, the standard complex structure on \mathbb{C}^n .
- ω is a **Kähler form** on \mathbb{C}^n and (\mathbb{C}^n, J_0) is a **Kähler manifold**.

- $\tau(z) = \pi \operatorname{Tr}[(-J_0^2(z))^{1/2}] = 2\pi n, z \in X.$
- The Kodaira-Laplacian:

$$\square^{L^p} = \bar{\partial}^{L^p*} \bar{\partial}^{L^p} = -\frac{1}{2} \sum_{j=1}^n \left(\frac{\partial}{\partial z_j} - \pi p \bar{z}_j \right) \frac{\partial}{\partial \bar{z}_j}.$$

- For the renormalized Bochner-Laplacian, we have

$$\Delta_p = 2\square^{L^p}.$$

$$\begin{aligned} \Delta_p &= - \sum_{j=1}^n \left[(\nabla_{\partial/\partial x_j}^{L^p})^2 + (\nabla_{\partial/\partial y_j}^{L^p})^2 \right] - 2\pi np \\ &= - \sum_{j=1}^n \left[\left(\frac{\partial}{\partial x_j} - \pi p \bar{z}_j \right)^2 + \left(\frac{\partial}{\partial y_j} - \pi i p \bar{z}_j \right)^2 \right] - 2\pi np \\ &= - \sum_{j=1}^n \left[\frac{\partial^2}{\partial x_j^2} + \frac{\partial^2}{\partial y_j^2} - \pi p \bar{z}_j \frac{\partial}{\partial \bar{z}_j} \right]. \end{aligned}$$

The almost complex structure

- $J_0 : TX \rightarrow TX$ a skew-adjoint linear endomorphism such that

$$\omega(u, v) = g(J_0 u, v), \quad u, v \in TX;$$

- $J : TX \rightarrow TX$ the linear endomorphism given by

$$J = J_0(-J_0^2)^{-1/2}.$$

- J is an almost complex structure on X , $J^2 = -Id_{TX}$, compatible with ω and g :

$$\omega(Ju, Jv) = \omega(u, v), \quad g(Ju, Jv) = g(u, v), \quad u, v \in TX.$$

- ω is positive: for $u \in TX \setminus 0$

$$\omega(u, Ju) = -g(JJ_0 u, u) = g((-J_0^2)^{1/2} u, u) > 0.$$

- If $J_0 = J$ and J is integrable, then (X, J) is a Kähler manifold.

Spectral gap property

Theorem (Guillemin-Urbe, 1988; Ma-Marinescu, 2002)

There exists $C_L > 0$ such that for any p

$$\sigma(\Delta_p) \subset [-C_L, C_L] \cup [2p\mu_0 - C_L, +\infty),$$

where the constant μ_0 is given by

$$\mu_0 = \inf_{u \in T_x X, x \in X} \frac{iR_x^L(u, J(x)u)}{|u|_g^2}.$$

Example

$$\mu_0 = \inf_{u \in TX} \frac{|(B^*B)^{1/4}u|_g^2}{|u|_g^2} = \inf_{x \in X} \inf (B^*B(x))^{1/2}.$$

Generalized Bergman projection

- \mathcal{H}_p the linear subspace of $L^2(X, L^p)$ spanned by the eigensections of Δ_p corresponding to eigenvalues in $[-C_L, C_L]$ (small eigenvalues).
- $P_{\mathcal{H}_p}$ the orthogonal projection in $L^2(X, L^p)$ onto \mathcal{H}_p (generalized Bergman projection).

Example

(X, ω) a compact Kähler manifold, L a holomorphic line bundle:

- $\Delta_p = 2\Box^{L^p}$, where \Box^{L^p} is the Kodaira Laplacian on L^p :

$$\sigma(\Delta_p) \subset \{0\} \cup [2p\mu_0 - C_L, +\infty),$$

- \mathcal{H}_p is the space $H^0(X, L^p)$ of holomorphic sections of L^p .
- $P_{\mathcal{H}_p}$ the usual Bergman projection.

Toeplitz operators

A **Toeplitz operator** is a sequence of bounded linear operators

$$T_p : L^2(X, L^p) \rightarrow L^2(X, L^p), p \in \mathbb{N}:$$

- For any $p \in \mathbb{N}$, we have

$$T_p = P_{\mathcal{H}_p} T_p P_{\mathcal{H}_p}.$$

- There exists a sequence $g_l \in C^\infty(X)$ such that

$$T_p = P_{\mathcal{H}_p} \left(\sum_{l=0}^{\infty} p^{-l} g_l \right) P_{\mathcal{H}_p} + \mathcal{O}(p^{-\infty}),$$

i.e. for any natural k there exists $C_k > 0$ such that

$$\left\| T_p - P_{\mathcal{H}_p} \left(\sum_{l=0}^k p^{-l} g_l \right) P_{\mathcal{H}_p} \right\| \leq C_k p^{-k-1}.$$

Toeplitz operators

For any p , the operator T_p acts on a finite-dimensional space \mathcal{H}_p . The dimension of \mathcal{H}_p is given, for p large enough, by the Riemann-Roch-Hirzebruch formula

$$d_p := \dim \mathcal{H}_p = \langle \text{ch}(L^p) \text{Td}(TX), [X] \rangle.$$

Here $\text{ch}(L^p)$ is the Chern character of L^p and $\text{Td}(TX)$ is the Todd class of the tangent bundle TX considered as a complex vector bundle with complex structure J . In particular,

$$d_p \sim p^n \int_X \frac{\omega^n}{n!}, \quad p \rightarrow \infty.$$

Algebra of Toeplitz operators

Theorem (Yu.K., 2017, Ios-Lo-Ma-Marinescu, 2017)

The product $T_{f,p}T_{g,p}$ of the Toeplitz operators

$$T_{f,p} = P_{\mathcal{H}_p} f P_{\mathcal{H}_p}, \quad T_{g,p} = P_{\mathcal{H}_p} g P_{\mathcal{H}_p}, \quad f, g \in C^\infty(X),$$

is a Toeplitz operator. It admits the asymptotic expansion

$$T_{f,p}T_{g,p} = \sum_{r=0}^{\infty} p^{-r} T_{C_r(f,g),p} + \mathcal{O}(p^{-\infty}),$$

with some $C_r(f, g) \in C^\infty(X)$, where C_r are bidifferential operators:

$$C_0(f, g) = fg, \quad C_1(f, g) - C_1(f, g) = i\{f, g\},$$

where $\{f, g\}$ is the Poisson bracket on $(X, 2\pi\omega)$.

Wells and localization of eigenfunctions

A self-adjoint Toeplitz operator T_ρ with principal symbol h :

$$T_\rho = P_{\mathcal{H}_\rho} \left(\sum_{l=0}^{\infty} \rho^{-l} g_l \right) P_{\mathcal{H}_\rho} + \mathcal{O}(\rho^{-\infty}), \quad g_0 = h.$$

An obvious lower bound:

$$(T_\rho v, v) \geq (\inf_{x \in X} h(x) + \mathcal{O}(\rho^{-1})) \|v\|^2, \quad v \in L^2(X, L^p).$$

Indeed, for $T_{h,\rho} = P_{\mathcal{H}_\rho} h P_{\mathcal{H}_\rho}$ and $v \in \mathcal{H}_\rho$,

$$(T_{h,\rho} u, u) = \int_X h(x) |u(x)|^2 dv_g(x) \geq \inf_{x \in X} h(x) \|u\|^2.$$

Wells and localization of eigenfunctions

Let $h_0 \in \mathbb{R}$. Suppose λ_p is a sequence of eigenvalues of T_p such that

$$\lambda_p \leq h_1 < h_0, \quad p \in \mathbb{N},$$

then the corresponding normalized eigensection u_p of T_p :

$$T_{h,p}u_p = \lambda_p u_p, \quad Pu_p = u_p, \quad \|u_p\| = 1$$

should be localized in the well

$$U_{h_0} = \{x \in X : h(x) \leq h_0\}.$$

Classically forbidden domain

$$X \setminus U_{h_0} = \{x \in X : h(x) > h_0\}.$$

Tunneling estimates

Assume that there exist $C > 0$ and $a > 0$ such that for any $p \in \mathbb{N}$ and $(x, y) \in X \times X$,

$$|(T_p - P_{\mathcal{H}_p} h P_{\mathcal{H}_p})(x, y)| < C p^{-1} e^{-a\sqrt{p}d(x,y)}.$$

Theorem (Y.K., 2018)

Suppose that $u_p \in \mathcal{H}_p$, $\|u_p\| = 1$, is a sequence of eigenfunctions of T_p , $T_p u_p = \lambda_p u_p$, such that

$$\lambda_p \leq h_1 < h_0, \quad p \in \mathbb{N}.$$

There exist $\alpha > 0$ and $C_1 > 0$ such that, for any $p \in \mathbb{N}$,

$$\int_X e^{2\alpha\sqrt{p}d(x, U_{h_0})} |u_p(x)|^2 dv_g(x) < C_1.$$

A self-adjoint Toeplitz operator T_p with principal symbol h :

$$T_p = P_{\mathcal{H}_p} \left(\sum_{l=0}^{\infty} p^{-l} g_l \right) P_{\mathcal{H}_p} + \mathcal{O}(p^{-\infty}), \quad g_0 = h.$$

Without loss of generality, we will assume that the principal symbol h satisfies the condition:

$$\min_{x \in X} h(x) = 0.$$

The spectrum of T_p consists of a finite number of eigenvalues

$$\lambda_p^0 \leq \lambda_p^1 \leq \dots \leq \lambda_p^{d_p-1}.$$

The asymptotic properties of λ_p^m in the semiclassical limit $p \rightarrow \infty$.

The Bergman projection

Let $x_0 \in X$ be a non-degenerate minimum of h :

$$\text{Hess } h(x_0) > 0.$$

The model operator for Δ_p at x_0 :

$$\mathcal{L}_0 = - \sum_{j=1}^{2n} \left(\frac{\partial}{\partial e_j} + \frac{1}{2} R_{x_0}^L(Z, e_j) \right)^2 - \tau(x_0),$$

- $\{e_j\}_{j=1, \dots, 2n}$ is an orthonormal base in $T_{x_0}X$.
- $\frac{\partial}{\partial U}$ the ordinary differentiation operator on $T_{x_0}X$ in the direction $U \in T_{x_0}X$.

\mathcal{P}_{x_0} the orthogonal projection in $L^2(T_{x_0}X)$ to the kernel of \mathcal{L}_0 (the Bergman projection of \mathcal{L}_0).

The Bergman kernel

$\mathcal{J}_{x_0} : T_{x_0}X \rightarrow T_{x_0}X$ is a skew-adjoint operator such that

$$R^L(u, v) = g(\mathcal{J}u, v), \quad u, v \in TX.$$

Actually, $\mathcal{J} = -2\pi iJ_0$.

We choose an orthonormal basis $\{e_j : j = 1, \dots, 2n\}$ of $T_{x_0}X$:

$$\mathcal{J}_{x_0} e_{2k-1} = a_k e_{2k}, \quad \mathcal{J}_{x_0} e_{2k} = -a_k e_{2k-1}, \quad a_k > 0 \quad k = 1, \dots, n.$$

We use this basis to define the coordinates Z on $T_{x_0}X \cong \mathbb{R}^{2n}$ as well as the complex coordinates z on $\mathbb{C}^n \cong \mathbb{R}^{2n}$, $z_j = Z_{2j-1} + iZ_{2j}$, $j = 1, \dots, n$.

The smooth Schwartz kernel (with respect to $dv_{TX}(Z)$) of \mathcal{P}_{x_0} :

$$\mathcal{P}(Z, Z') = \frac{1}{(2\pi)^n} \prod_{j=1}^n a_j \exp \left(-\frac{1}{4} \sum_j a_j (|z_j|^2 + |z'_j|^2 - 2z_j \bar{z}'_j) \right).$$

The Toeplitz operator \mathcal{T}_{x_0} in $L^2(T_{x_0}X)$ (the model operator for T_p at x_0):

$$\mathcal{T}_{x_0} = \mathcal{P}_{x_0}(q_{x_0}(Z) + g_1(x_0))\mathcal{P}_{x_0},$$

- $q_{x_0}(Z)$ the second order term in the Taylor expansion of h at x_0 (in normal coordinates near x_0):

$$q_{x_0}(Z) = \left(\frac{1}{2} \text{Hess } h(x_0) Z, Z \right), \quad Z \in T_{x_0}X,$$

ia positive quadratic form on $T_{x_0}X \cong \mathbb{R}^{2n}$.

- g_1 is the second coefficient in the asymptotic expansion for $\{T_p\}$.

It is an unbounded self-adjoint operator in $L^2(T_{x_0}X)$ with discrete spectrum. The eigenvalues of \mathcal{T}_{x_0} do not depend on the choice of normal coordinates, and the lowest eigenvalue is simple.

Eigenvalue asymptotic expansions

Assume that the principal symbol h satisfies the following conditions:

- $h(x) \geq 0$ for any $x \in X$;
- $\min_{x \in X} h(x) = 0$;
- Each minimum is non-degenerate.

Then $U_0 = \{x \in X : h(x) = 0\}$ is a finite set (discrete wells):

$$U_0 = \{x_1, \dots, x_N\}.$$

Let \mathcal{T} be the self-adjoint operator on $L^2(T_{x_1}X) \oplus \dots \oplus L^2(T_{x_N}X)$

$$\mathcal{T} = \mathcal{T}_{x_1} \oplus \dots \oplus \mathcal{T}_{x_N}.$$

Theorem (Y.K. (2018), A. Deleporte (2017, for Kähler manifolds))

Let $\{\lambda_p^m\}$ be the increasing sequence of the eigenvalues of T_p on \mathcal{H}_p (counted with multiplicities) and $\{\mu_m\}$ the increasing sequence of the eigenvalues of \mathcal{T} (counted with multiplicities). Then, for any fixed m , λ_p^m has an asymptotic expansion, when $p \rightarrow \infty$, of the form

$$\lambda_p^m = p^{-1} \mu_m + p^{-3/2} \phi_m + \mathcal{O}(p^{-2}).$$

Theorem (Y.K. (2018), A. Deleporte (2017, for Kähler manifolds))

If μ is a simple eigenvalue of \mathcal{T}_{x_j} for some j , then there exists a sequence λ_p of eigenvalues of T_p on \mathcal{H}_p which has a complete asymptotic expansion of the form

$$\lambda_p^0 \sim p^{-1} \sum_{k=0}^{+\infty} a_k p^{-k}, \quad a_0 = \mu.$$

Toeplitz operator $\mathcal{T}(Q)$ in $L^2(\mathbb{C}^n)$

$$\mathcal{T}(Q) = \mathcal{P}Q : \ker \mathcal{L}_0 \subset L^2(\mathbb{C}^n) \rightarrow \ker \mathcal{L}_0 \subset L^2(\mathbb{C}^n),$$

where $Q = Q(z, \bar{z})$ is a polynomial in \mathbb{C}^n and \mathcal{P} is the orthogonal projection in $L^2(\mathbb{C}^n)$ to the kernel of \mathcal{L}_0 .

Fock space \mathcal{F}_n is the space of holomorphic functions F in \mathbb{C}^n such that $e^{-\frac{1}{2}|z|^2} F \in L^2(\mathbb{C}^n)$.

\mathcal{F}_n is a closed subspace in $L^2(\mathbb{C}^n; e^{-\frac{1}{2}|z|^2} dz)$.

The orthogonal projection $\Pi : L^2(\mathbb{C}^n; e^{-\frac{1}{2}|z|^2} dz) \rightarrow \mathcal{F}_n$:

$$\Pi F(z) = \frac{1}{\pi^n} \int_{\mathbb{C}^n} \exp\left(-|z'|^2 + z \cdot \bar{z}'\right) F(z') dz' d\bar{z}'.$$

Consider the isometry $S : L^2(\mathbb{C}^n; e^{-\frac{1}{2}|z|^2} dz) \rightarrow L^2(\mathbb{C}^n)$ given, for $u \in L^2(\mathbb{C}^n; e^{-\frac{1}{2}|z|^2} dz)$, by

$$Su(z) = \frac{\prod_{j=1}^n a_j}{2^n} e^{-\frac{1}{4} \sum_j a_j |z_j|^2} u(\phi(z)),$$

where $\phi : \mathbb{C}^n \rightarrow \mathbb{C}^n$ is a linear isomorphism given by

$$\phi(z) = \left(\frac{\sqrt{a_1}}{\sqrt{2}} z_1, \dots, \frac{\sqrt{a_n}}{\sqrt{2}} z_n \right), \quad z \in \mathbb{C}^n.$$

It is easy to see that $S\Pi S^{-1} = \mathcal{P}$. It follows that $S(\mathcal{F}_n) = \ker \mathcal{L}_0$ and

$$\mathcal{T}(Q) = ST^0(Q \circ \phi^{-1})S^{-1}.$$

where $\mathcal{T}^0(Q)$ is a Toeplitz operator in the Fock space defined by

$$\mathcal{T}^0(Q) = \Pi Q : \mathcal{F}_n \rightarrow \mathcal{F}_n.$$

Bargmann transform

The Bargmann transform is an isometry $B : L^2(\mathbb{R}^n) \rightarrow \mathcal{F}_n$ defined by

$$Bf(z) = \pi^{-n/4} \int_{\mathbb{R}^n} \exp \left[- \left(\frac{1}{2} z \cdot z + \frac{1}{2} x \cdot x - \sqrt{2} z \cdot x \right) \right] f(x) dx, \quad z \in \mathbb{C}^n.$$

For the position and momentum operators in the Schrödinger presentation $\hat{q}_k = x_k$, $\hat{p}_k = \frac{1}{i} \frac{\partial}{\partial x_k}$, the corresponding operators in the Bargmann-Fock presentation

$$B\hat{q}_k B^{-1} = \frac{1}{\sqrt{2}} \left(z_k + \frac{\partial}{\partial z_k} \right), \quad B\hat{p}_k B^{-1} = \frac{1}{\sqrt{2}} i \left(z_k - \frac{\partial}{\partial z_k} \right).$$

For the creation and annihilation operators $\hat{a}_k = \frac{1}{\sqrt{2}}(\hat{q}_k + i\hat{p}_k)$, $\hat{a}_k^* = \frac{1}{\sqrt{2}}(\hat{q}_k - i\hat{p}_k)$, the corresponding operators in the Bargmann-Fock presentation

$$B\hat{a}_k B^{-1} = \frac{\partial}{\partial z_k}, \quad B\hat{a}_k^* B^{-1} = z_k.$$

- For any $F \in L^2(\mathbb{C}^n; e^{-\frac{1}{2}|z|^2} dz)$, $\frac{\partial}{\partial z_k} \Pi F = \Pi \bar{z}_k F$;
- For any $F \in \mathcal{F}_n$, $z_k \Pi F = \Pi(z_k F)$.

Let P be a polynomial in \mathbb{C}^n . If we write P as

$$P(\bar{z}, z) = \sum_{k,l} A_{k;l} \bar{z}_1^{k_1} \dots \bar{z}_n^{k_n} z_1^{l_1} \dots z_n^{l_n},$$

then, for any $F \in \mathcal{F}_n$,

$$\mathcal{T}^0(P)F = \Pi(P(\bar{z}, z)F) = P(\partial_z, z)F,$$

where $P(\partial_z, z)$ is the operator in \mathcal{F}_n given by

$$P(\partial_z, z) = \sum_{k,l} A_{k;l} \frac{\partial^{k_1}}{\partial z_1^{k_1}} \dots \frac{\partial^{k_n}}{\partial z_n^{k_n}} z_1^{l_1} \dots z_n^{l_n}.$$

Anti-Wick symbols (Berezin, 1971)

- Under the Bargmann transform B , the operator $\mathcal{T}^0(P)$ in \mathcal{F}_n corresponds to the operator $B^{-1}\mathcal{T}^0(P)B$ in $L^2(\mathbb{R}^n)$ given by

$$B^{-1}\mathcal{T}^0(P)B = \sum A_{k;l} \hat{a}^{k_1} \dots \hat{a}^{k_n} (\hat{a}_1^*)^{l_1} \dots (\hat{a}_n^*)^{l_n}.$$

- The operator $B^{-1}\mathcal{T}^0(P)B$ is the differential operator with polynomial coefficients in \mathbb{R}^n ;
- $P(\bar{z}, z)$ is the anti-Wick symbol of $B^{-1}\mathcal{T}^0(P)B$ (the Schrödinger presentation) and $\mathcal{T}^0(P)$ (the Bargmann-Fock presentation);
- Some sufficient conditions of self-adjointness of the operator $B^{-1}P(\partial_z, z)B$ (Berezin, 1971).

One can compute the Weyl symbol of this operator by the well-known formula (Berezin, 1971). If P is a positive definite quadratic form, then

$$B^{-1}\mathcal{T}^0(P)B = Op_w(\tilde{P}) + \frac{\text{tr}(\tilde{P})}{2},$$

where \tilde{P} is a quadratic form on \mathbb{R}^{2n} , corresponding to P under the linear isomorphism $(x, \xi) \in \mathbb{R}^{2n} \cong T^*\mathbb{R}^n \mapsto z \in \mathbb{C}^n$:

$$z_k = \frac{1}{\sqrt{2}}(x_k - i\xi_k), \quad k = 1, \dots, n.$$

$Op_w(\tilde{P})$ is the pseudodifferential operator in \mathbb{R}^n with Weyl symbol \tilde{P} .

The Bochner-Laplacian Δ^{L^p} :

$$\Delta^{L^p} = (\nabla^{L^p})^* \nabla^{L^p} : C^\infty(X, L^p) \rightarrow C^\infty(X, L^p).$$

τ is a smooth function on X :

$$\tau(x) = \pi \operatorname{Tr}[(-\mathcal{J}_0^2(x))^{1/2}], \quad x \in X.$$

For the magnetic Laplacian $\Delta^L = -\sum_{j=1}^d \left(\frac{\partial}{\partial z_j} - iA_j(z) \right)^2$

$$\tau(x) = \frac{1}{2} \operatorname{Tr}(B^* B)^{1/2} = \operatorname{Tr}^+(B).$$

Assumptions

- $\min_{x \in X} \tau(x) = \tau_0$;
- There exists a unique $x_0 \in X$ such that $\tau(x) = \tau_0$, which is non-degenerate:

$$\operatorname{Hess} \tau(x_0) > 0.$$

Schrödinger operator type representation

$$\Delta^{L^p} = (\nabla^{L^p})^* \nabla^{L^p} = \Delta_p + p\tau,$$

the renormalized Bochner-Laplacian Δ_p satisfies the gap property:

$$\sigma(\Delta_p) \subset [-C_L, C_L] \cup [2p\mu_0 - C_L, +\infty),$$

\mathcal{H}_p corresponds to the lowest Landau levels.

Upper estimates for $\lambda_j(\Delta^{L^p})$ (the Rayleigh-Ritz technique):

$$\lambda_j(\Delta^{L^p}) \leq \lambda_j(P_{\mathcal{H}_p} \Delta^{L^p} P_{\mathcal{H}_p}), \quad j \in \mathbb{N}.$$

The operator

$$\rho^{-1} P_{\mathcal{H}_\rho} \Delta^{L\rho} P_{\mathcal{H}_\rho} = \rho^{-1} P_{\mathcal{H}_\rho} \Delta_\rho P_{\mathcal{H}_\rho} + P_{\mathcal{H}_\rho} \tau P_{\mathcal{H}_\rho}$$

is a Toeplitz operator:

$$\rho^{-1} P_{\mathcal{H}_\rho} \Delta^{L\rho} P_{\mathcal{H}_\rho} = P_{\mathcal{H}_\rho} \left(\sum_{l=0}^{\infty} \rho^{-l} g_l \right) P_{\mathcal{H}_\rho} + \mathcal{O}(\rho^{-\infty}).$$

The leading term

$$g_0(x) = \tau(x).$$

The next term is the principal symbol of $\Delta_\rho P_{\mathcal{H}_\rho}$:

$$g_1(x) = J_{1,2}(x).$$

Computation of $J_{1,2}$

Put

$$\frac{\partial}{\partial z_j} = \frac{1}{2} \left(\frac{\partial}{\partial Z_{2j-1}} - i \frac{\partial}{\partial Z_{2j}} \right), \quad \frac{\partial}{\partial \bar{z}_j} = \frac{1}{2} \left(\frac{\partial}{\partial Z_{2j-1}} + i \frac{\partial}{\partial Z_{2j}} \right).$$

Let $\mathcal{R}(Z) = \sum_{j=1}^{2n} Z_j e_j = Z$ denote the radial vector field on $T_{x_0} X$.Define first order differential operators $b_j, b_j^+, j = 1, \dots, n$, on $T_{x_0} X$ by

$$b_j = -2\nabla_{\frac{\partial}{\partial z_j}} - R_{x_0}^L(\mathcal{R}, \frac{\partial}{\partial z_j}) \quad b_j^+ = 2\nabla_{\frac{\partial}{\partial \bar{z}_j}} + R_{x_0}^L(\mathcal{R}, \frac{\partial}{\partial \bar{z}_j}).$$

So we can write

$$\mathcal{L}_{x_0} = - \sum_{j=1}^{2n} \left(\frac{\partial}{\partial e_j} + \frac{1}{2} R_{x_0}^L(Z, e_j) \right)^2 - \tau(x_0) = \sum_{j=1}^n b_j b_j^+.$$

$$J_{1,2}(x_0) = \frac{F_{1,2,x_0}(0,0)}{\mathcal{P}_{x_0}(0,0)}, \quad F_{1,2,x_0}(Z, Z') = [\mathcal{P}_{x_0} \mathcal{F}_{1,2,x_0} \mathcal{P}_{x_0}](Z, Z'),$$

where $\mathcal{F}_{1,2,x_0}$ is an unbounded linear operator in $L^2(T_{x_0}X)$ given by

$$\begin{aligned} \mathcal{F}_{1,2,x_0} = & 4 \left\langle R_{x_0}^{TX} \left(\frac{\partial}{\partial z_i}, \frac{\partial}{\partial z_j} \right) \frac{\partial}{\partial \bar{z}_i}, \frac{\partial}{\partial \bar{z}_j} \right\rangle \\ & + \left\langle (\nabla^X \nabla^X \mathcal{J})_{(\mathcal{R}, \mathcal{R})} \frac{\partial}{\partial z_i}, \frac{\partial}{\partial \bar{z}_i} \right\rangle + \frac{\sqrt{-1}}{4} \operatorname{tr}_{|TX} \left(\nabla^X \nabla^X (\mathcal{J}\mathcal{J}) \right)_{(\mathcal{R}, \mathcal{R})} \\ & + \frac{1}{9} |(\nabla_{\mathcal{R}}^X \mathcal{J})\mathcal{R}|^2 + \frac{4}{9} \left\langle (\nabla_{\mathcal{R}}^X \mathcal{J})\mathcal{R}, \frac{\partial}{\partial z_i} \right\rangle b_i^+ \mathcal{L}_0^{-1} b_i \left\langle (\nabla_{\mathcal{R}}^X \mathcal{J})\mathcal{R}, \frac{\partial}{\partial \bar{z}_i} \right\rangle. \end{aligned}$$

If $J_0 = J$ (almost-Kähler), then

$$J_{1,2}(x_0) = \frac{1}{24} |\nabla^X J|_{x_0}^2.$$

Here if $\{e_j\}$ is a local orthonormal frame of (TX, g^{TX}) , then

$$|\nabla^X J|^2 = \sum_{ij} |(\nabla_{e_i}^X J)e_j|^2.$$

Upper bounds for eigenvalues

Consider the Toeplitz operator \mathcal{T}_{x_0} in $L^2(T_{x_0}X)$ defined by

$$\mathcal{T}_{x_0} = \mathcal{P}_{x_0}(q_{x_0}(Z) + \mathcal{J}_{1,2}(x_0))\mathcal{P}_{x_0}.$$

where

$$q_{x_0}(Z) = \left(\frac{1}{2} \text{Hess } \tau(x_0) Z, Z \right), \quad Z \in T_{x_0}X.$$

Let $\{\lambda_j(\Delta^{L^p})\}$ be the increasing sequence of the eigenvalues of the operator Δ^{L^p} (counted with multiplicities) and $\{\mu_j\}$ be the increasing sequence of the eigenvalues of \mathcal{T}_{x_0} (counted with multiplicities).

Theorem (Yu. K. (2018))

For any $j \in \mathbb{N}$, there exists $\phi_j \in \mathbb{R}$ such that

$$\lambda_j(\Delta^{L^p}) \leq p\tau_0 + \mu_j + p^{-1/2}\phi_j + \mathcal{O}(p^{-1}), \quad p \rightarrow \infty.$$

2D-magnetic Laplacian

In particular, for the 2D-magnetic Laplacian on a Riemann surface X :

$$\mathbf{B} = b(x)dv_g, \quad dv_g = \sqrt{g}dx_1 \wedge dx_2,$$

assume $b(x) > 0$ for any $x \in X$ and there exists a unique x_0 such that

$$b(x_0) = b_0 := \min_{x \in X} b(x),$$

if we denote

$$a = \text{Tr} \left(\frac{1}{2} \text{Hess } b(x_0) \right)^{1/2}, \quad d = \det \left(\frac{1}{2} \text{Hess } b(x_0) \right).$$

we have (Helffer-Y.K., 2010, Helffer-Morame, 2001):

$$\lambda_j(\Delta^{L^p}) \leq pb_0 + \left[\frac{2d^{1/2}}{b_0} j + \frac{a^2}{2b_0} \right] + C_j p^{-1/2}, \quad j \in \mathbb{N}.$$